A Compact, Totally Passive, Multi-Pass Slab Laser Amplifier Based on Stable, Degenerate Optical Resonators

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Why do we need compact multipass amplifiers?

- Availability of small, high repetition rate, picosecond pulse sources (e.g microchip and SESAM oscillators) and photon-counting techniques is opening up a wide range of new applications, e.g.
 - compact photon-counting SLR stations (SLR2000)
 - Airborne/spaceborne 3D imaging lidars
 - Lunar and interplanetary transponders
- Most of these applications require pulse energies between 40 and 4000 μ J whereas the oscillators produce sub- μ J (10 ps SESAM) up to 250 μ J (custom 400 psec microchip).
- At several KHz repetition rates with CW diode pumping of the amplifiers, the gain per pass can be relatively low so highly multi-passed amplifiers are desirable and more efficient.

Unsaturated Multipass Gain

G



 $G_0 = single pass unsaturated gain$ $G_{mp} = G_0^{2N} = multipass unsaturated gain$

Multipass Amplifier Approaches

Passive Amplifier/Multiple Mirrors (e.g. Q-Peak laser in SLR2000)





Stable Degenerate Resonators*

- A <u>stable</u> optical resonator is defined by two spherical mirrors with radii of curvature, b_1 and b_2 , and separated by a distance d such that $0 \le (1-d/b_1)(1-d/b_2) \le 1$
- At certain mirror separations, d, the resonator becomes "degenerate" and can be characterized by an integer N
- At a mirror separation with "degeneracy" N:
 - The Hermite-Gaussian resonator modes divide into N discrete frequencies separated by c/2NL where L is the resonator length; thus, N=1 represents the highest degeneracy where all spatial modes oscillate at the same frequency.
 - Hole-coupled lasers exhibit large power losses because the frequency-degenerate modes can couple together to create a low loss composite mode with a null at the coupling hole
 - Ray paths can be defined which repeat themselves after N round trips (useful for multipass amplifiers)

*Reference: I. A. Ramsay and J. J. Degnan, "A Ray Analysis of Optical Resonators Formed by Two Spherical Mirrors", Applied Optics, Vol. 9, pp. 385-398, February, 1970.

General Resonator

If b_1 and b_2 are the mirror radii of curvature, the degenerate mirror separations are given by

$$d_{\pm}(N,K) = \frac{b_1 + b_2}{2} \pm \frac{1}{2}\sqrt{b_1^2 + b_2^2 + 2b_1b_2\cos\left(\frac{2\pi K}{N}\right)}$$

where N is the degeneracy factor, K = 0 for N = 1, and $1 \le K \le \frac{N}{2}$ for N > 1 provided K >1 is not divisible into N N = 1 2 3 4 5 6 7 8 K= 0 1 1 1 1 1 1 1 1 2 2 3 2 3

*Reference: I. A. Ramsay and J. J. Degnan, "A Ray Analysis of Optical Resonators Formed by Two Spherical Mirrors", Applied Optics, Vol. 9, pp. 385-398, February, 1970.

Symmetric Resonator $(b_1 = b_2 = b)$

$$d_{\pm}(N,K) = b \left[1 \pm \cos\left(\frac{\pi K}{N}\right) \right]$$



Symmetric Resonator $(b_1 = b_2)$ N=4, K=1

Ecliptic

Non-Ecliptic





Flat-Concave Resonator $(b_2 = \infty)$



Resonator Length (normalized to b1)

Flat-Concave Resonator ($b_2 = \infty$)

Non-Ecliptic

Ecliptic

a а t t **Ecliptic Ray Trace** Non-Ecliptic Ray Trace s i \mathbf{S} 10 20 i D D 1 $\mathbf{N} = \mathbf{4}$ 1 а а 0 0 i i $\mathbf{K} = \mathbf{1}$ d d а а R R -10-200.2 0.8 0.2 0.8 0 0.4 0.6 0 0.4 0.6 Normalized Resonator Distance Normalized Resonator Distance a a t t Ecliptic Ray Trace Non-Ecliptic Ray Trace \mathbf{S} s 40 20 i i D D 20 N = 51 a а i i C $\mathbf{K} = \mathbf{2}$ d d a -20 R а R -20 -40 0.2 0.4 0.8 0.8 0 0.6 0 0.2 0.4 0.6 1 Normalized Resonator Distance Normalized Resonator Distance

Ecliptic vs Non-Ecliptic Amplifiers

Ecliptic

•Amplifier entrance and exit rays overlap

•Requires optical isolation from the oscillator

•Requires a means to insert and extract the beam (e.g. polarization rotation)

•Samples less of the laser slab (less efficient extraction?)

•Generally easier to align since the insertion axis always lies along mirror normal

•With one flat mirror, a variable pass amplifier can be constructed using one translatable mirror.

Non-Ecliptic •Amplifier entrance and exit rays do not

overlap

•Probably does not require optical isolation from the oscillator

•Does not require a means to insert and extract the beam (e.g. polarization rotation)

•Samples more of the laser slab (more efficient extraction?)

•Generally more difficult to align since the insertion angle varies with degeneracy N.

Ecliptic Variable Pass Amplifier



Excellent Beam Control: Collimated beam in, same size collimated beam out!

Summary

- Microchip and SESAM oscillators can generate picosecond pulses at multi-KHz rates but at low single pulse energies (microjoules or less)
- Many airborne and spaceborne applications require amplifications of 10 to 10³ in a compact, efficient, diode-pumped package.
- Since CW-diode pumped amps typically have low single pass gains, many passes through the amplifier may be required to reach the required pulse energies and to extract the stored energy efficiently
- Degenerate resonator multipass amplifiers can provide :
 - high multipass gain in a compact, easily aligned package
 - A fair amount of isolation from the oscillator and reduced internal feedback for suppressing self-oscillations
 - Variable number of passes with one translating mirror which can be set for optimum performance or compensate for a degradation in oscillator power
 - Excellent beam control since it preserves the gaussian parameters of the input beam at the output due to periodic refocusing